

# S-Duality and Chern-Simons Theory

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String Theory Seminar

University of Texas at Austin

## Based on

- Yoon Pyo Hong and OG, “S-duality and Chern-Simons Theory,” [[arXiv:hep-th/0812.1213](#)]
- Yoon Pyo Hong and OG, “S-twisted compactification of  $N = 4$ , Topological 2+1D Quantum Field Theory, and Minimal Strings” [[arXiv:hep-th/0902.????](#)]

# S-duality

$$\tau \equiv \frac{4\pi i}{g_{\text{YM}}^2} + \frac{\theta}{2\pi}$$

$$\mathbf{s} = \begin{pmatrix} \mathbf{a} & \mathbf{b} \\ \mathbf{c} & \mathbf{d} \end{pmatrix} \in \text{SL}(2, \mathbb{Z})$$

$$\tau \rightarrow \frac{\mathbf{a}\tau + \mathbf{b}}{\mathbf{c}\tau + \mathbf{d}}.$$

# S-duality's action on states

Temporal gauge:  $A_0 = 0$ .

$$\tilde{\Psi}(A) \equiv \int [\mathcal{D}\tilde{A}] \mathcal{S}(A, \tilde{A}) \Psi(\tilde{A})$$

$$\tau \rightarrow \frac{\mathbf{a}\tau + \mathbf{b}}{\mathbf{c}\tau + \mathbf{d}}, \quad E_i \rightarrow \mathbf{a}E_i + \mathbf{b}B_i, \quad B_i \rightarrow \mathbf{c}E_i + \mathbf{d}B_i.$$

[Lozano; Gaiotto & Witten]

$$\mathcal{S}(A, \tilde{A}) = \exp \left\{ \frac{i}{4\pi\mathbf{c}} \int (\mathbf{d}A \wedge dA - 2\tilde{A} \wedge dA + \mathbf{a}\tilde{A} \wedge d\tilde{A}) \right\}.$$

$$\tilde{E}_i \mathcal{S} = \mathcal{S}(\mathbf{a}E_i + \mathbf{b}B_i), \quad \tilde{B}_i \mathcal{S} = \mathcal{S}(\mathbf{c}E_i + \mathbf{d}B_i).$$

$$E_i \equiv -2\pi i \delta / \delta A_i$$

## $U(1)$ Chern-Simons from S-duality

$$\tilde{\Psi}\{A\} \equiv \int [\mathcal{D}\tilde{A}] \mathcal{S}(A, \tilde{A}) \Psi(\tilde{A})$$

$$\mathcal{S}(A, \tilde{A}) = \exp \left\{ \frac{i}{4\pi\mathbf{c}} \int (\mathbf{d}A \wedge dA - 2\tilde{A} \wedge dA + \mathbf{a}\tilde{A} \wedge d\tilde{A}) \right\}.$$

$$A = \tilde{A} \implies \mathcal{I}(A) \equiv \frac{\mathbf{a} + \mathbf{d} - 2}{4\pi\mathbf{c}} \int A \wedge dA.$$

$$\text{CS level: } k \equiv (\mathbf{a} + \mathbf{d} - 2)/\mathbf{c}.$$

Physical interpretation?

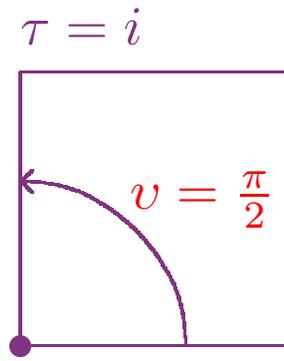
# Selfduality

$$\tau = \frac{\mathbf{a}\tau + \mathbf{b}}{\mathbf{c}\tau + \mathbf{d}} \implies \mathbf{c}\tau + \mathbf{d} = e^{i\nu}.$$

At a selfdual  $\tau$  we can compactify on a circle with an S-twist.

$$\mathbf{s} \equiv \begin{pmatrix} \mathbf{a} & \mathbf{b} \\ \mathbf{c} & \mathbf{d} \end{pmatrix} = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$$

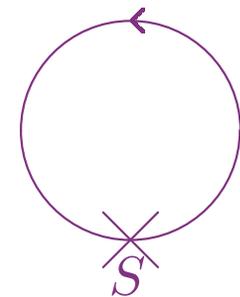
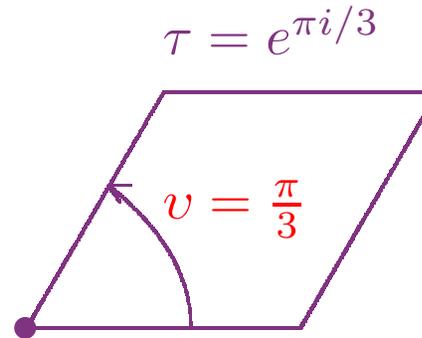
$$\tau \rightarrow -\frac{1}{\tau} \quad |k| = \left| \frac{\mathbf{a}+\mathbf{d}-2}{\mathbf{c}} \right| = 2$$



periodic time

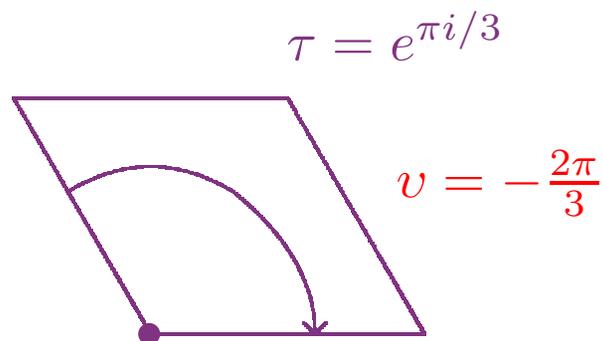
$$\mathbf{s} \equiv \begin{pmatrix} \mathbf{a} & \mathbf{b} \\ \mathbf{c} & \mathbf{d} \end{pmatrix} = \begin{pmatrix} 1 & -1 \\ 1 & 0 \end{pmatrix}$$

$$\tau \rightarrow \frac{\tau-1}{\tau} \quad |k| = \left| \frac{\mathbf{a}+\mathbf{d}-2}{\mathbf{c}} \right| = 1$$



$$\mathbf{s} = \begin{pmatrix} -1 & 1 \\ -1 & 0 \end{pmatrix}$$

$$\tau \rightarrow \frac{\tau-1}{\tau} \quad |k| = 3$$



# $N = 4$ Super Yang-Mills

$A_\mu$	gauge field	$\mu = 0 \dots 3$
$\Phi^I$	adjoint-valued scalars	$I = 1 \dots 6$
$\psi_\alpha^a$	adjoint-valued spinors	$a = 1 \dots 4$ and $\alpha = 1, 2$
$\bar{\psi}_{a\dot{\alpha}}$	complex conjugate spinors	$a = 1 \dots 4$ and $\dot{\alpha} = \dot{1}, \dot{2}$
$Q_{a\alpha}$	SUSY generators	$a = 1 \dots 4$ and $\alpha = 1, 2$
$\bar{Q}_{\dot{\alpha}}^a$	complex conjugate generators	$a = 1 \dots 4$ and $\dot{\alpha} = \dot{1}, \dot{2}$

$$Z^1 = \Phi^1 + i\Phi^4, \quad Z^2 = \Phi^2 + i\Phi^5, \quad Z^3 = \Phi^3 + i\Phi^6.$$

# Supersymmetry

$$\mathbf{s} : \tau \rightarrow \frac{\mathbf{a}\tau + \mathbf{b}}{\mathbf{c}\tau + \mathbf{d}}$$

$$\mathbf{s} : Q_{a\alpha} \rightarrow \left( \frac{\mathbf{c}\tau + \mathbf{d}}{|\mathbf{c}\tau + \mathbf{d}|} \right)^{1/2} Q_{a\alpha} = e^{\frac{i\nu}{2}} Q_{a\alpha}$$

[Kapustin & Witten]

$$\mathbf{s} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \implies \nu = \frac{\pi}{2}$$

# R-Symmetry

$$Spin(6) \simeq SU(4)$$

$$\gamma \equiv \begin{pmatrix} e^{i\varphi_1} & & & \\ & e^{i\varphi_2} & & \\ & & e^{i\varphi_3} & \\ & & & e^{i\varphi_4} \end{pmatrix} \in SU(4), \quad \left( \sum_a \varphi_a = 0 \right),$$

acts as

$$\gamma(\psi_\alpha^a) = e^{i\varphi_a} \psi_\alpha^a, \quad \gamma(\bar{\psi}_{a\alpha}) = e^{-i\varphi_a} \bar{\psi}_{a\alpha}, \quad a = 1 \dots 4.$$

$$\gamma(Z^k) = e^{i(\varphi_k + \varphi_4)} Z^k, \quad k = 1 \dots 3.$$

## Combined R-S- action

$$Q_{a\alpha} \rightarrow e^{\frac{iv}{2} - i\varphi_a} Q_{a\alpha}.$$

$\implies N = 2r$  invariant generators

$$r = \#\{a \text{ for which } e^{i\varphi_a} = e^{iv/2}\}$$

# R- and S- twisted boundary conditions



$$\Phi(x = 0^-) = \gamma[\Phi(x = 0^+)]$$

$$Z^k(x = 0^-) = e^{i(\varphi_k + \varphi_4)} \Phi(x = 0^+), \quad k = 1, 2, 3$$

...



$$\Psi(A, \dots)|_{t=0^+} = \int [\mathcal{D}\tilde{A}] \mathcal{S}(A, \tilde{A}) \Psi(\tilde{A}, \dots)|_{t=0^-}$$

# SUSY in 2+1D

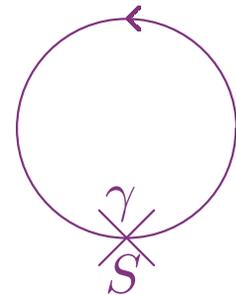
$$\implies N = 2r, \quad r = \#\{a \text{ for which } e^{i\varphi_a} = e^{iv/2}\}$$

$$\gamma = \begin{pmatrix} e^{\frac{i}{2}v} & & & \\ & e^{\frac{i}{2}v} & & \\ & & e^{\frac{i}{2}v} & \\ & & & e^{-\frac{3i}{2}v} \end{pmatrix} \implies N = 6$$

$$\gamma = \begin{pmatrix} e^{\frac{i}{2}v} & & & \\ & e^{\frac{i}{2}v} & & \\ & & e^{-i(v+\varphi_4)} & \\ & & & e^{i\varphi_4} \end{pmatrix} \implies N = 4$$

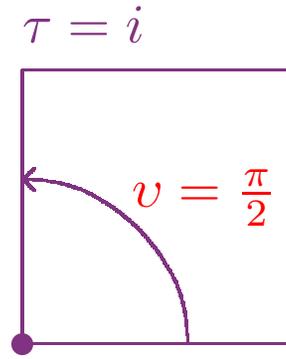
$\gamma =$  R-symmetry twist

$$e^{iv} \equiv \mathbf{c}\tau + \mathbf{d}$$



$$\mathbf{s} \equiv \begin{pmatrix} \mathbf{a} & \mathbf{b} \\ \mathbf{c} & \mathbf{d} \end{pmatrix} = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$$

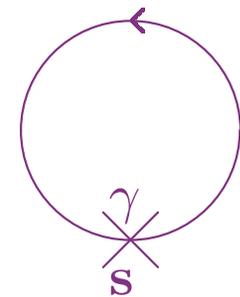
$$\tau \rightarrow -\frac{1}{\tau}$$



$$\mathbf{s}(\mathbf{a}, \mathbf{b}, \mathbf{c}, \mathbf{d})$$

$$\gamma(v)$$

$N = 4$  SYM

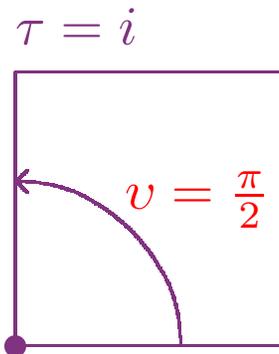


$N = 6$   
in 2+1D

IR???

$$\mathbf{s} \equiv \begin{pmatrix} \mathbf{a} & \mathbf{b} \\ \mathbf{c} & \mathbf{d} \end{pmatrix} = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$$

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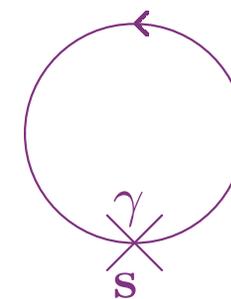
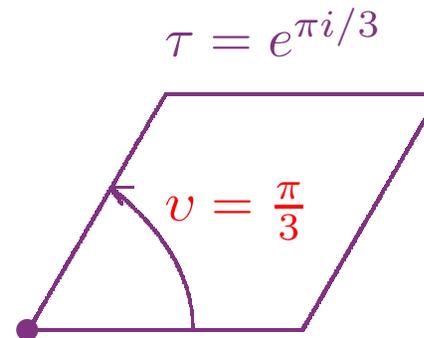


$\mathbf{s}(\mathbf{a}, \mathbf{b}, \mathbf{c}, \mathbf{d})$   
 $\gamma(v)$

$N = 4$  SYM

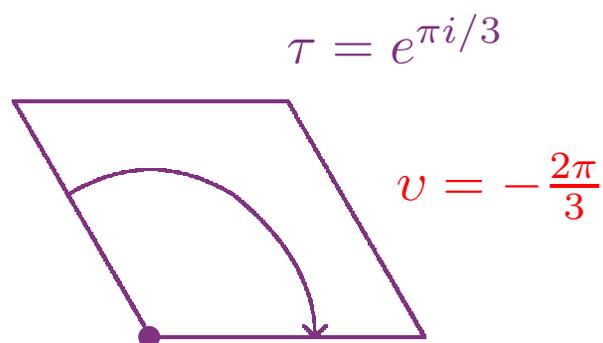
$$\mathbf{s} \equiv \begin{pmatrix} \mathbf{a} & \mathbf{b} \\ \mathbf{c} & \mathbf{d} \end{pmatrix} = \begin{pmatrix} 1 & -1 \\ 1 & 0 \end{pmatrix}$$

$$\tau \rightarrow \frac{\tau-1}{\tau}$$



$$\mathbf{s} = \begin{pmatrix} -1 & 1 \\ -1 & 0 \end{pmatrix}$$

$$\tau \rightarrow \frac{\tau-1}{\tau}$$



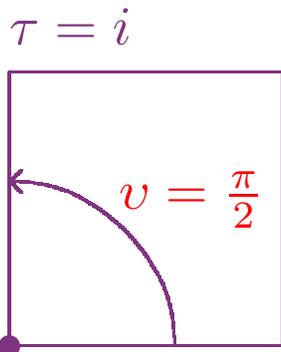
$N = 6$   
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IR???

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CS at  $k = 2$ ?



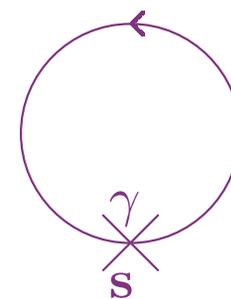
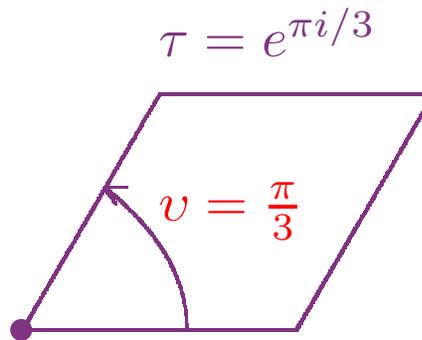
$\mathbf{s}(\mathbf{a}, \mathbf{b}, \mathbf{c}, \mathbf{d})$   
 $\gamma(\mathbf{v})$

$N = 4$  SYM

$$\mathbf{s} \equiv \begin{pmatrix} \mathbf{a} & \mathbf{b} \\ \mathbf{c} & \mathbf{d} \end{pmatrix} = \begin{pmatrix} 1 & -1 \\ 1 & 0 \end{pmatrix}$$

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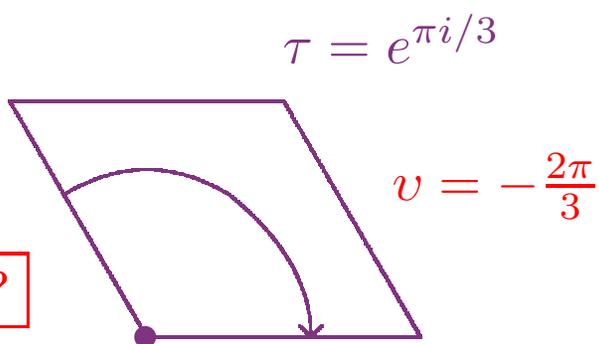
CS at  $k = 1$ ?



$$\mathbf{s} = \begin{pmatrix} -1 & 1 \\ -1 & 0 \end{pmatrix}$$

$$\tau \rightarrow \frac{\tau-1}{\tau}$$

CS at  $k = 3$ ?



$N = 6$   
in 2+1D

IR???

# Moduli

$$Z \equiv Z^1 \equiv \phi^1 + i\phi^4$$

BPS operators:

$$\mathcal{O}_p \equiv g_{\text{YM}}^{-p} \text{tr}(Z^p), \quad p = 1, 2, \dots$$

These operators are  $\text{SL}(2, \mathbb{Z})$ -duality invariant [Intriligator].

Action of R-symmetry twist:

$$(\mathcal{O}_p)^\gamma = e^{ipv} \mathcal{O}_p.$$

$\mathcal{O}_p$  is single-valued if and only if  $e^{ipv} = 1$ .

## Moduli . . .

- for  $\tau = i$  and  $\mathbf{s} = \begin{pmatrix} 1 & -1 \\ 1 & 0 \end{pmatrix}$ ,  $\langle \mathcal{O}_p \rangle \neq 0$  requires  $p \in 4\mathbb{Z}$ ;
- for  $\tau = e^{\pi i/3}$  and  $\mathbf{s} = \begin{pmatrix} 1 & -1 \\ 1 & 0 \end{pmatrix}$   $\langle \mathcal{O}_p \rangle \neq 0$  requires  $p \in 6\mathbb{Z}$ ;
- for  $\tau = e^{\pi i/3}$  and  $\mathbf{s} = \begin{pmatrix} -1 & 1 \\ -1 & 0 \end{pmatrix}$   $\langle \mathcal{O}_p \rangle \neq 0$  requires  $p \in 3\mathbb{Z}$ .

For  $U(n)$ ,  $\mathcal{O}_{n+1}, \mathcal{O}_{n+2}, \dots$  are not independent of  $\mathcal{O}_1, \dots, \mathcal{O}_n$ . Thus for  $\tau = i$  and  $\mathbf{s} = \mathbf{s}'$ , for example, if  $n < 4$  none of the operators  $\mathcal{O}_p$  can get a VEV.

## States on $T^2$ from String Theory

type	brane	1	2	3	4	5	6	7	8	9	10	
IIB	D3	=	=	×								T on 1:
IIA	D2	○	=	×								to M:
M	M2	○	=	×							○	on 2:
IIA	F1	○		×							○	

Legend:

- | direction doesn't exist in the theory;
- = a direction that the brane wraps;
- × a direction that the brane wraps and has the S-R-twist;
- a compact direction that the brane doesn't wrap;

# Counting fixed-points

type	brane	1	2	3	4	5	6	7	8	9	10
IIB	D3	=	=	×							
IIA	F1	○		×							○

$$\tau = i \implies g_{\text{IIB}} = 1 \implies R_1 = R_{10}.$$

Directions 1, 10 form a  $T^2$  of complex structure  $\tau$ ;

**F1-strings** are  $n$  points in directions 1, 10;

**F1-strings** are wound in direction 3;

## Counting fixed-points ...

Directions 1, 10 form a  $T^2$  of complex structure  $\tau$ ;

F1-strings are  $n$  points in directions 1, 10;

F1-strings are wound in direction 3;

S-R-twist is entirely geometrical!

It is a rotation by  $\nu = \pi/2$  of  $T^2$ ;

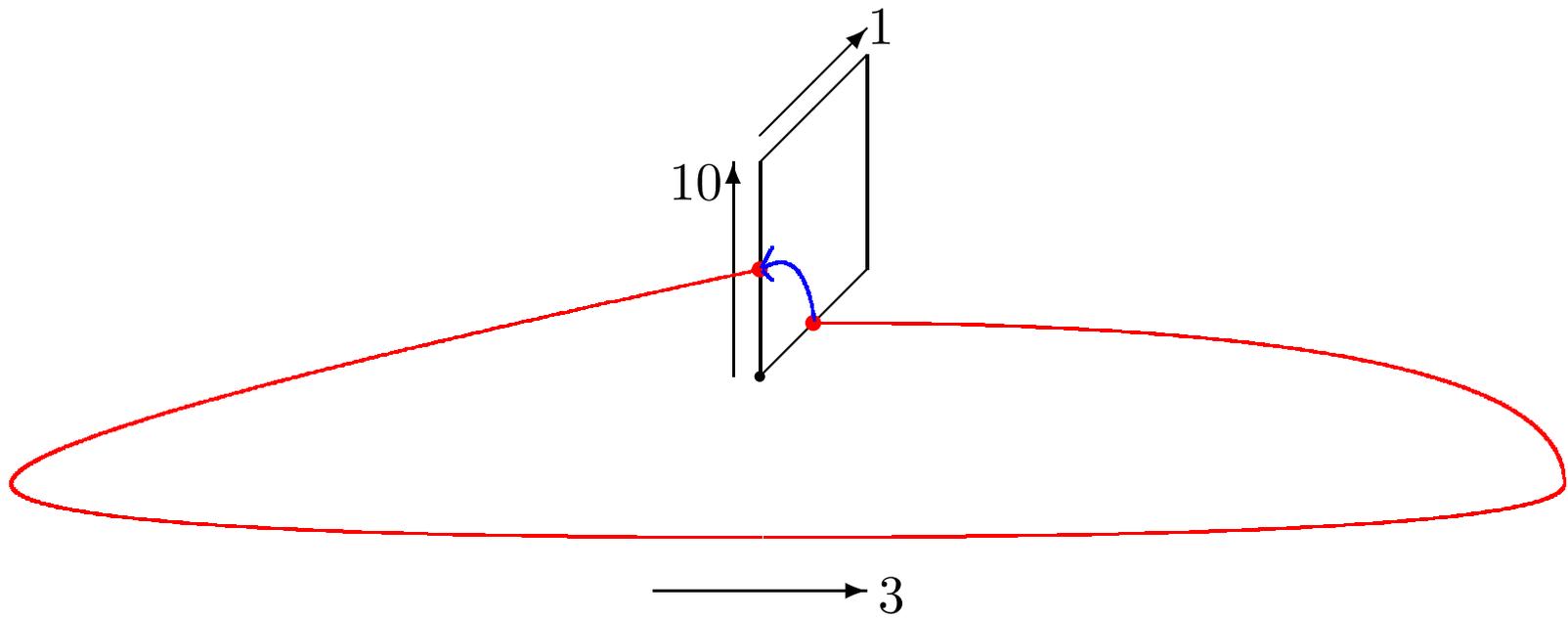
Need to find fixed points of this rotation (up to  $S_n$ );

$\{z_{\sigma(1)}, \dots, z_{\sigma(n)}\} = \{z_1, \dots, z_n\}$  up to  $\mathbb{Z} + \mathbb{Z}\tau$ ;

One Ramond-Ramond ground state for each fixed point.

$T^2$  (directions 1, 10) fibered over  $S^1$  (direction 3):

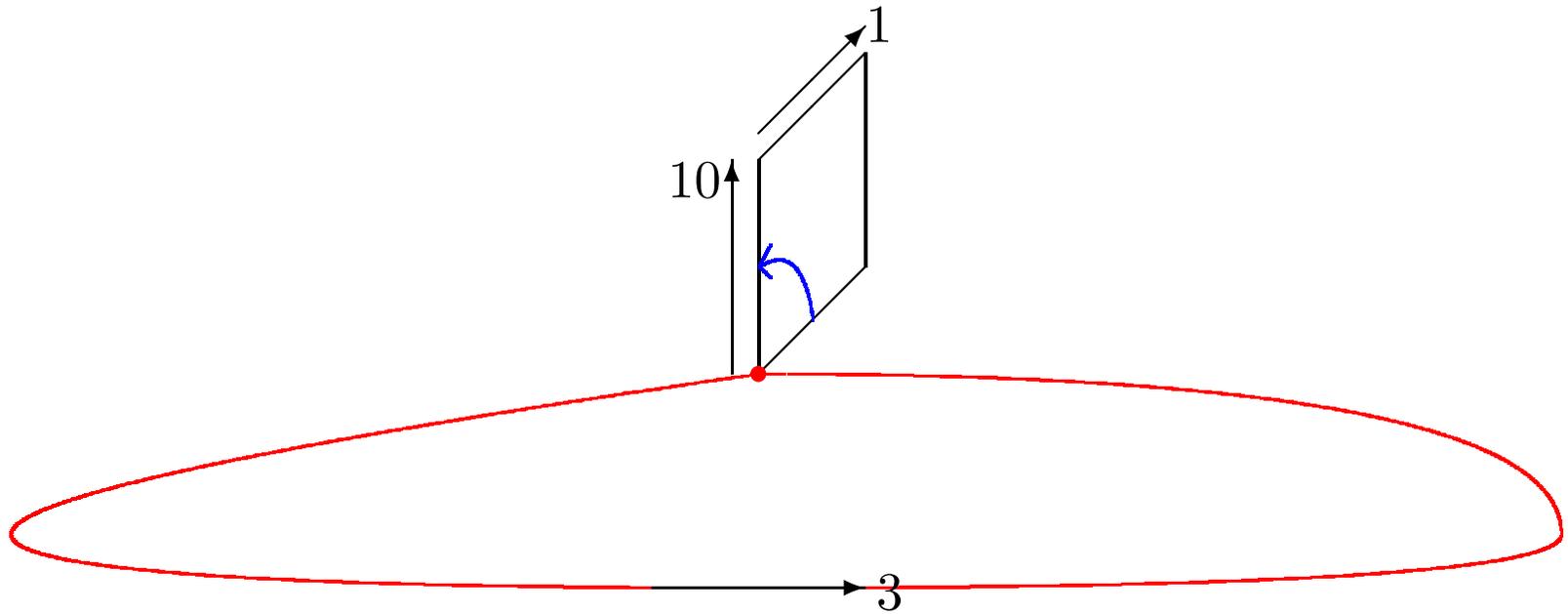
Geometrical twist and wound string



$T^2$  (directions 1, 10) fibered over  $S^1$  (direction 3):

Geometrical twist and wound string

Minimal energy configuration: find fixed points of twist!

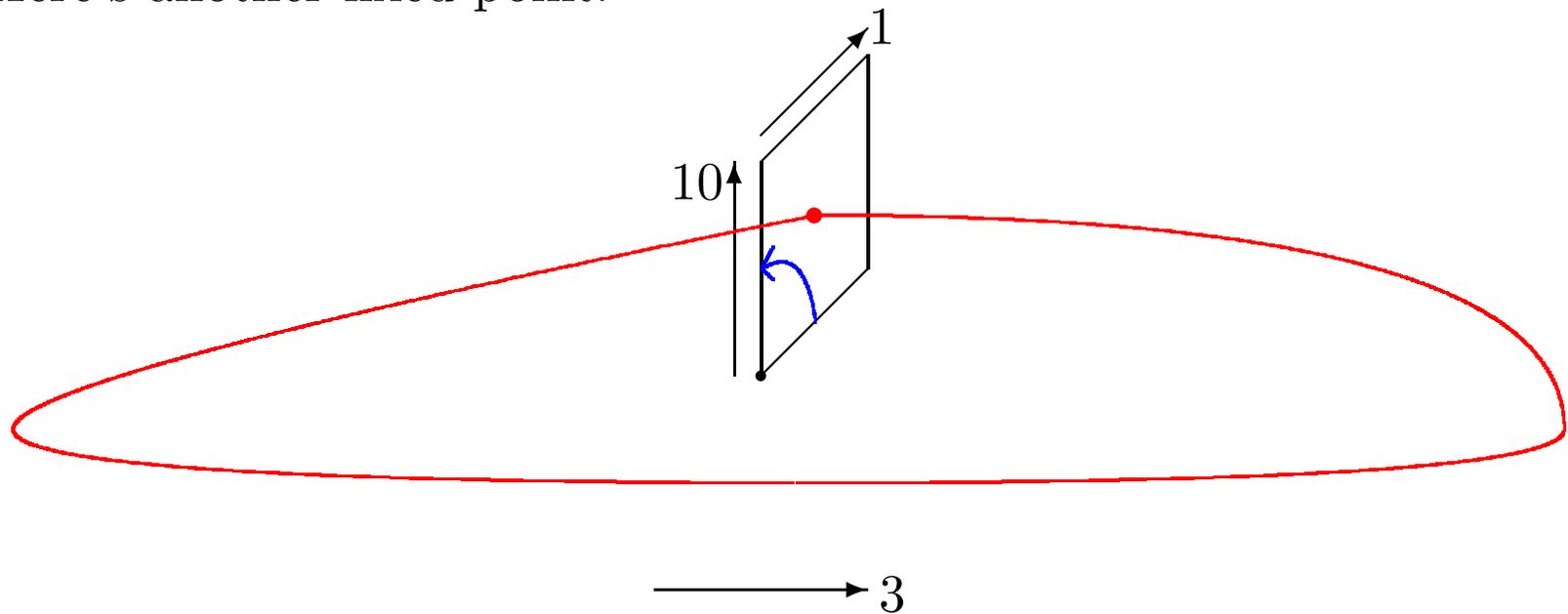


$T^2$  (directions 1, 10) fibered over  $S^1$  (direction 3):

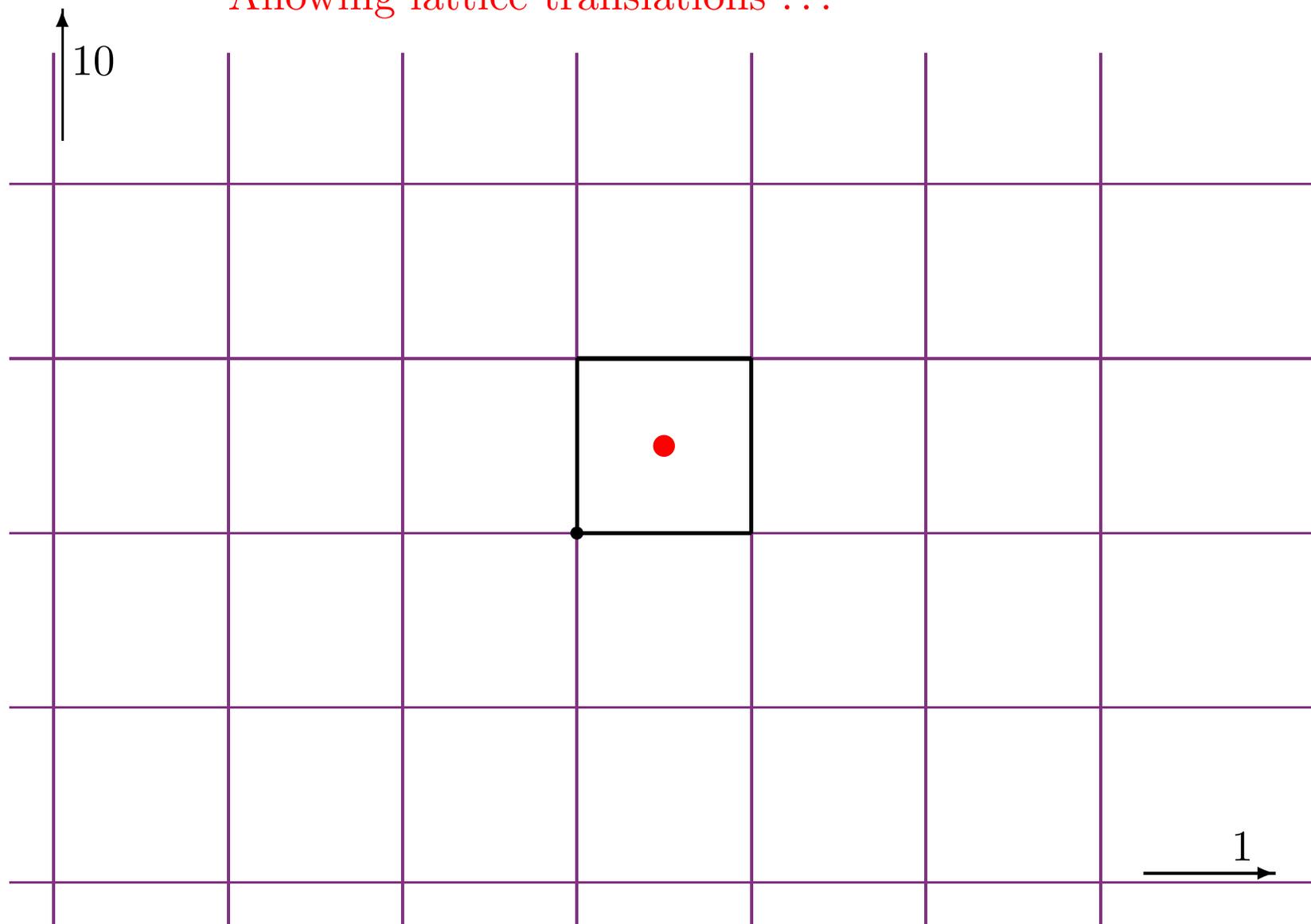
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Minimal energy configuration: find fixed points of twist!

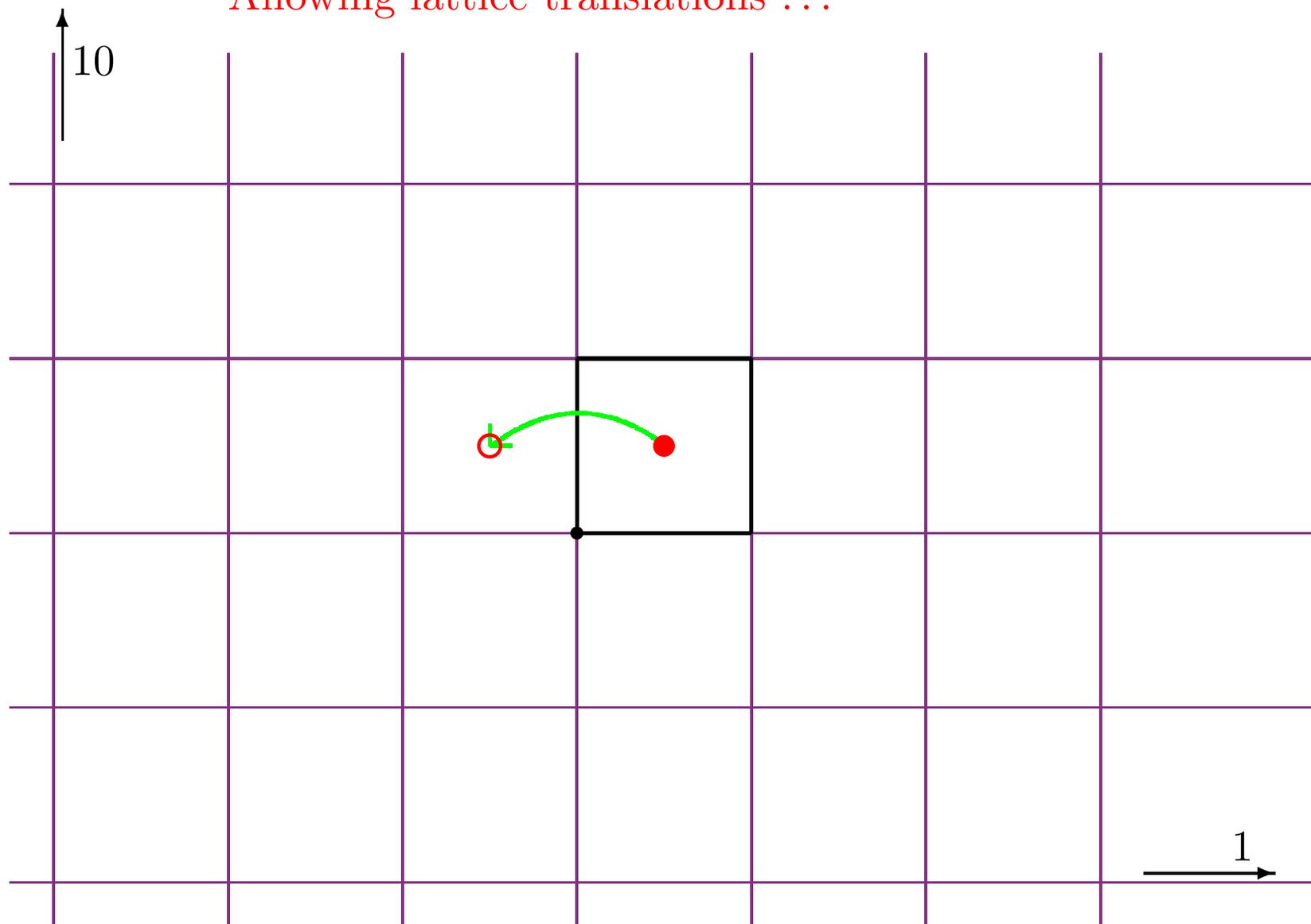
Here's another fixed point.



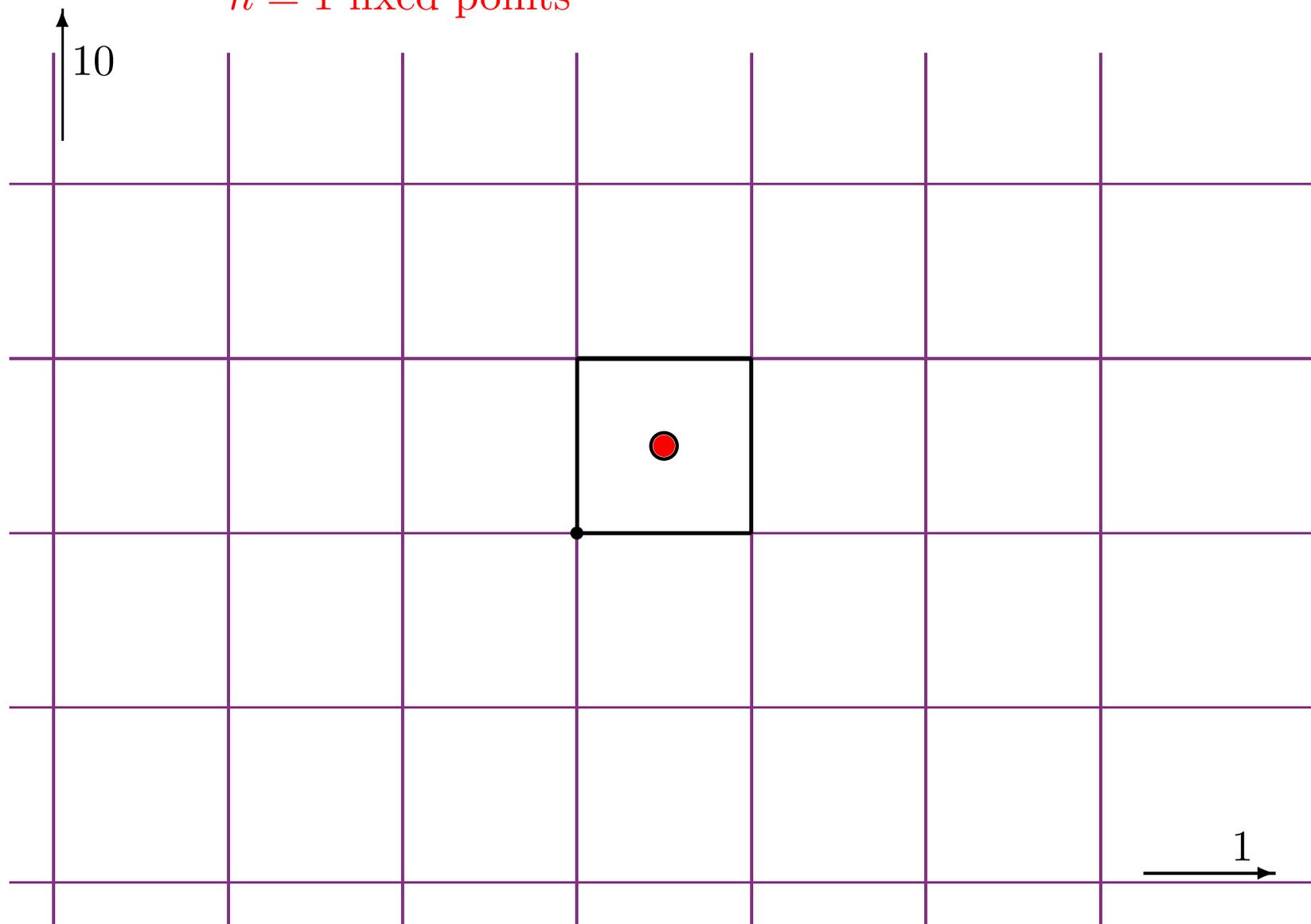
Allowing lattice translations ...



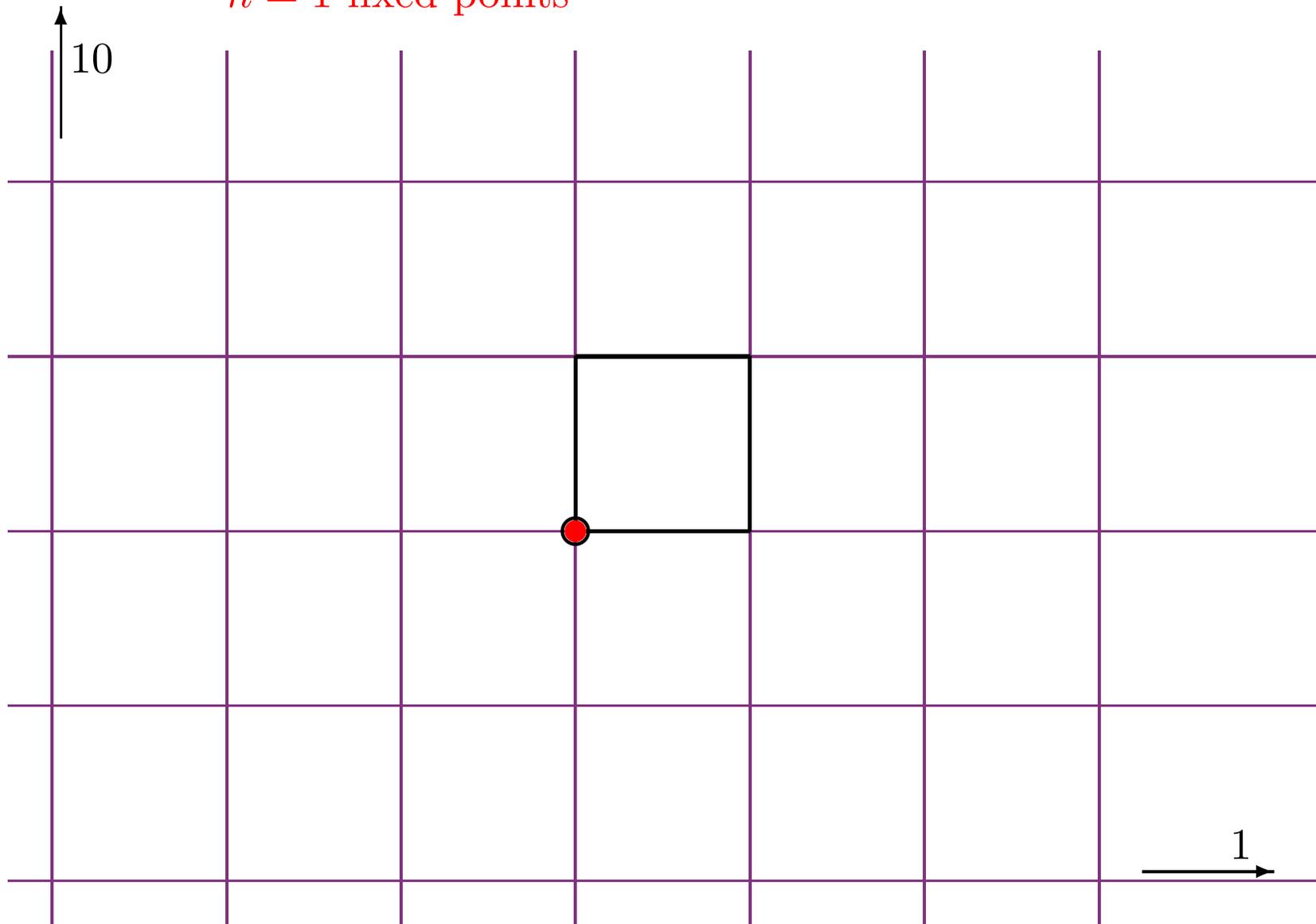
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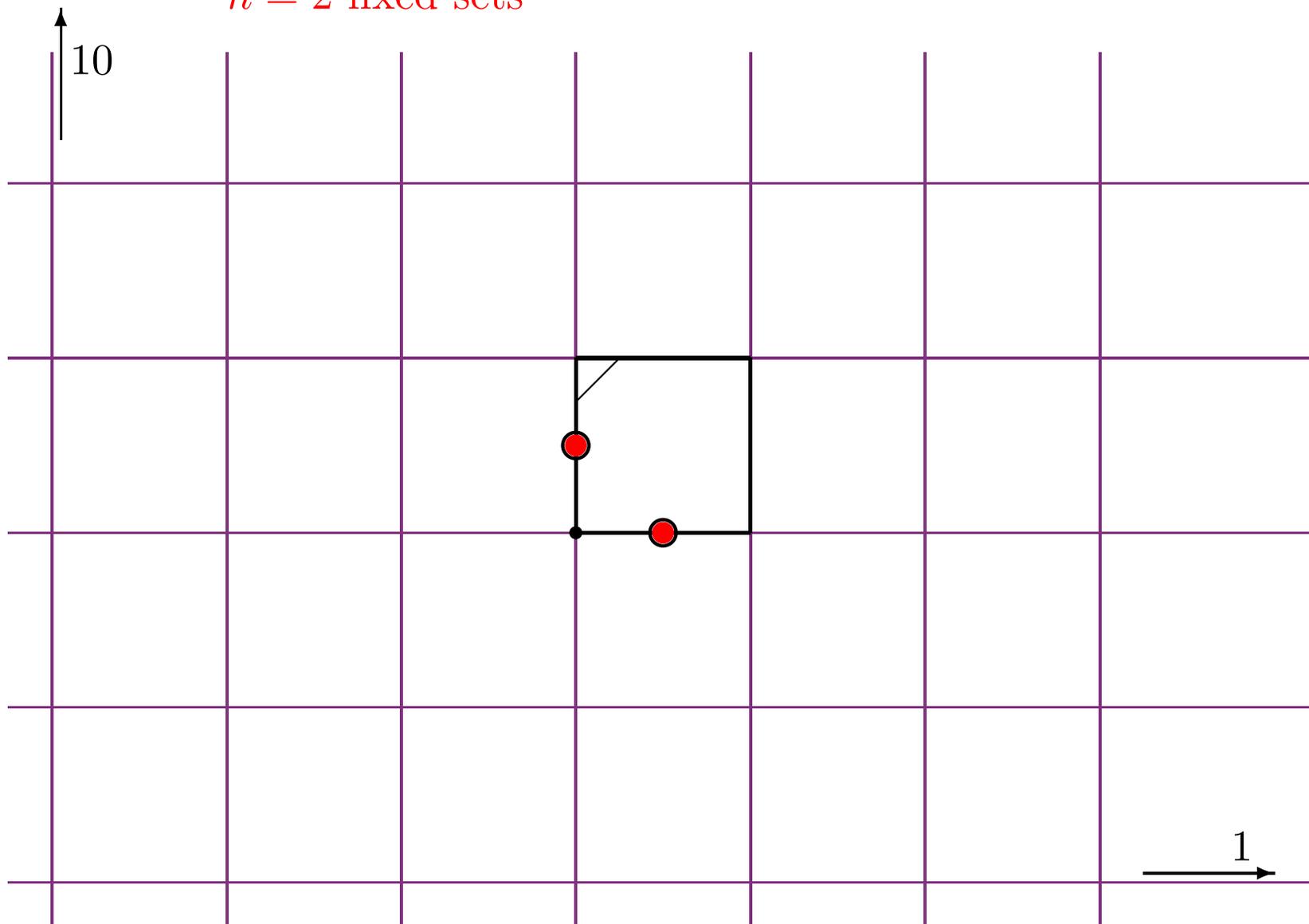
$n = 1$  fixed points



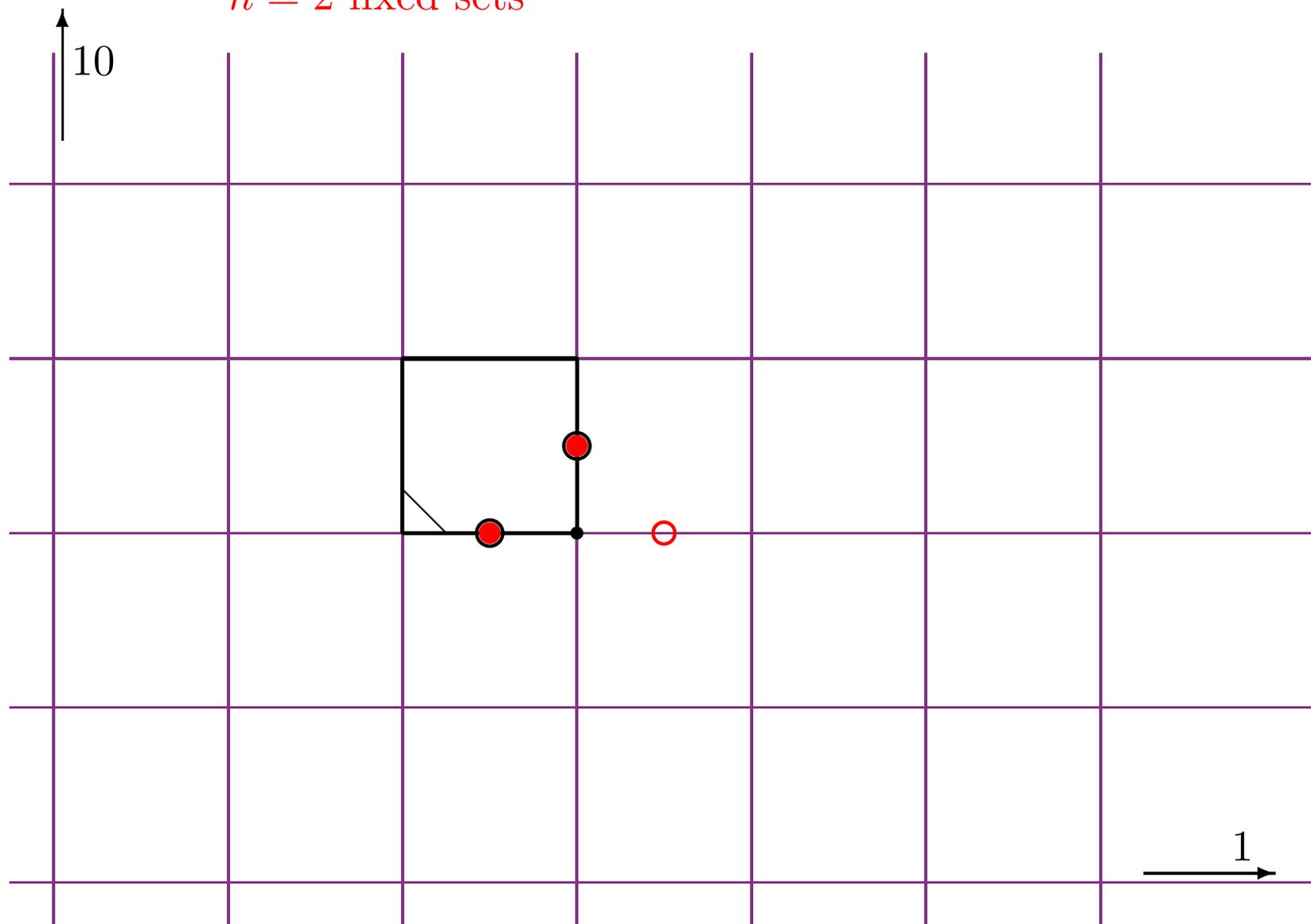
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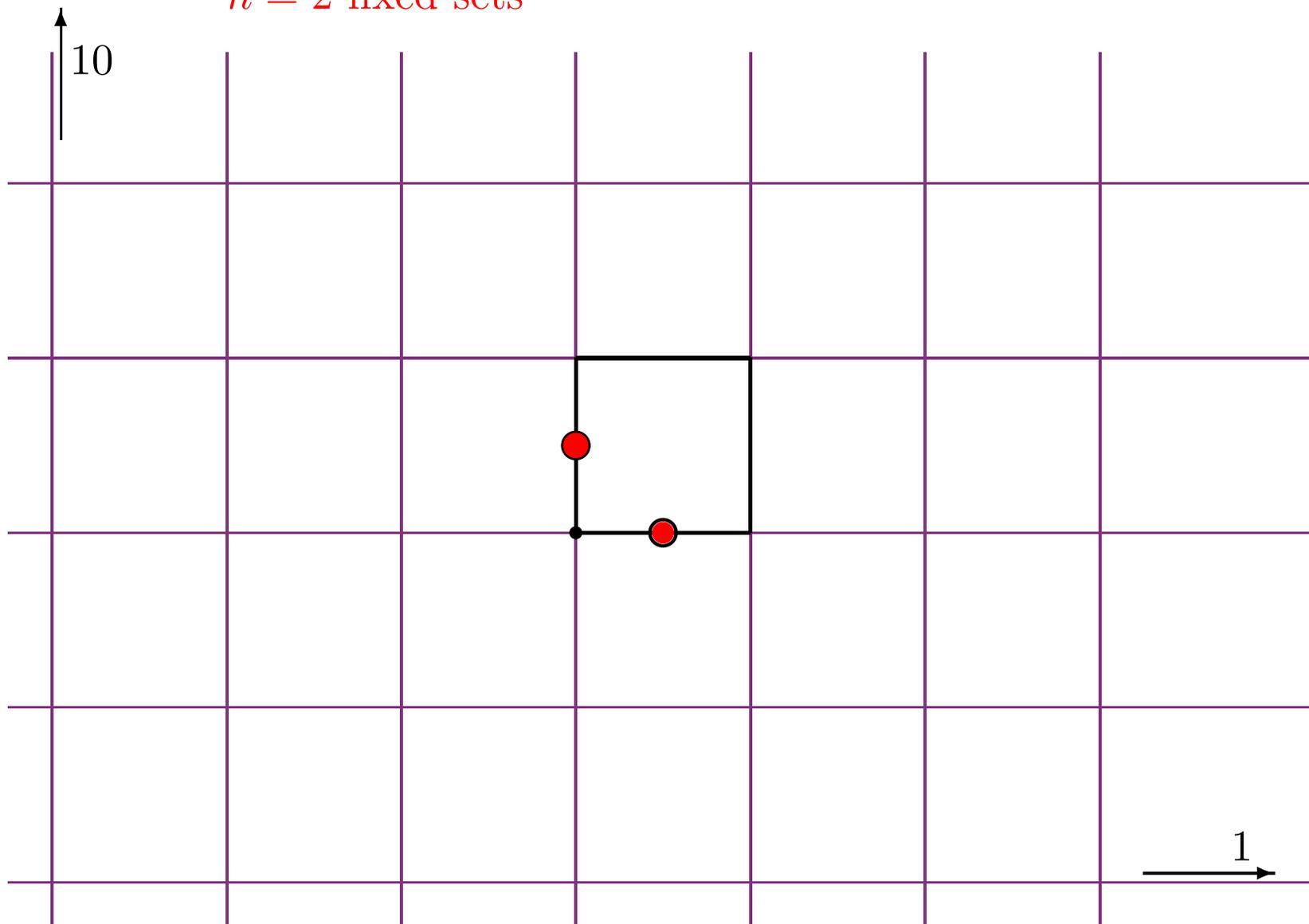
$n = 2$  fixed sets



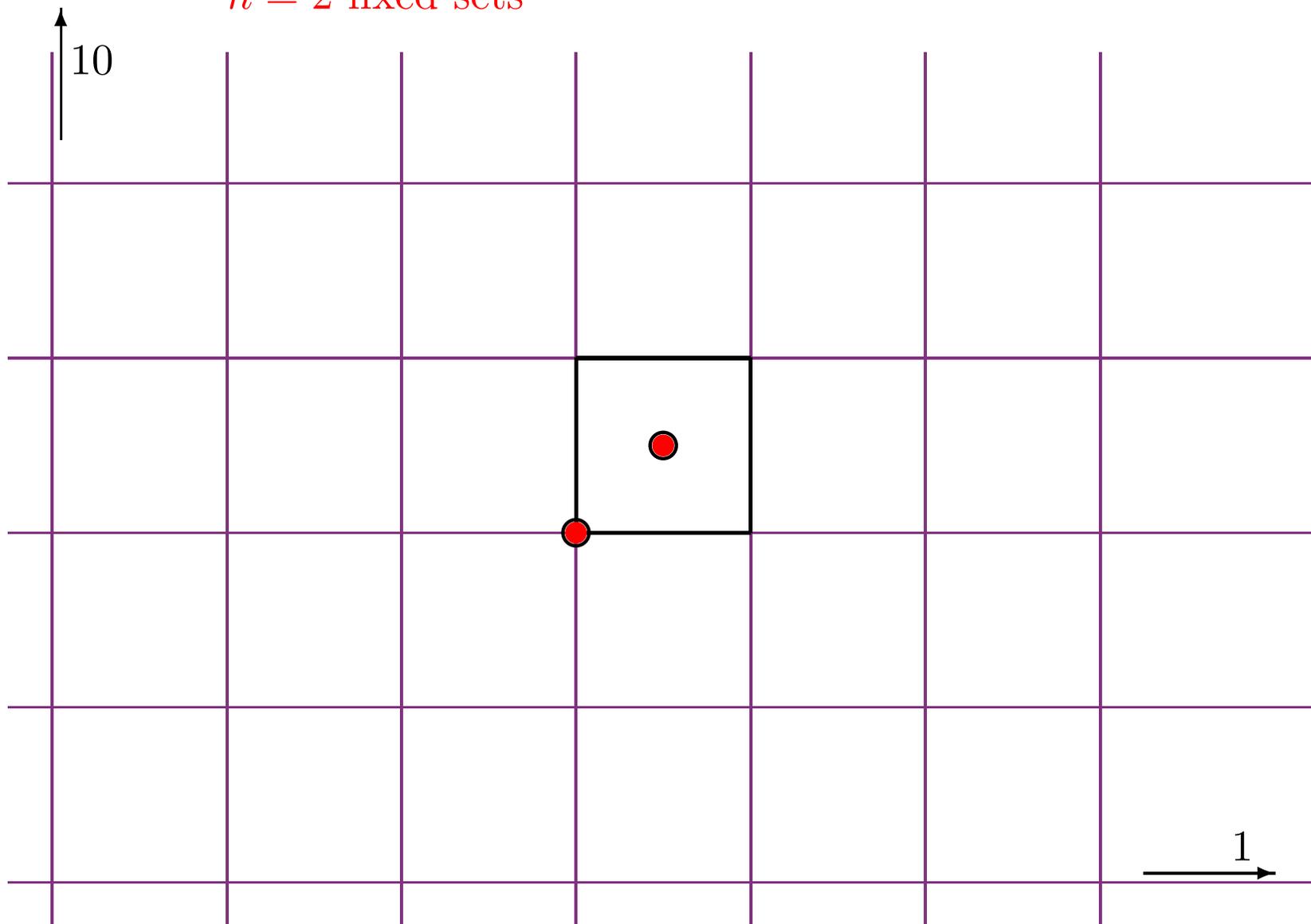
$n = 2$  fixed sets



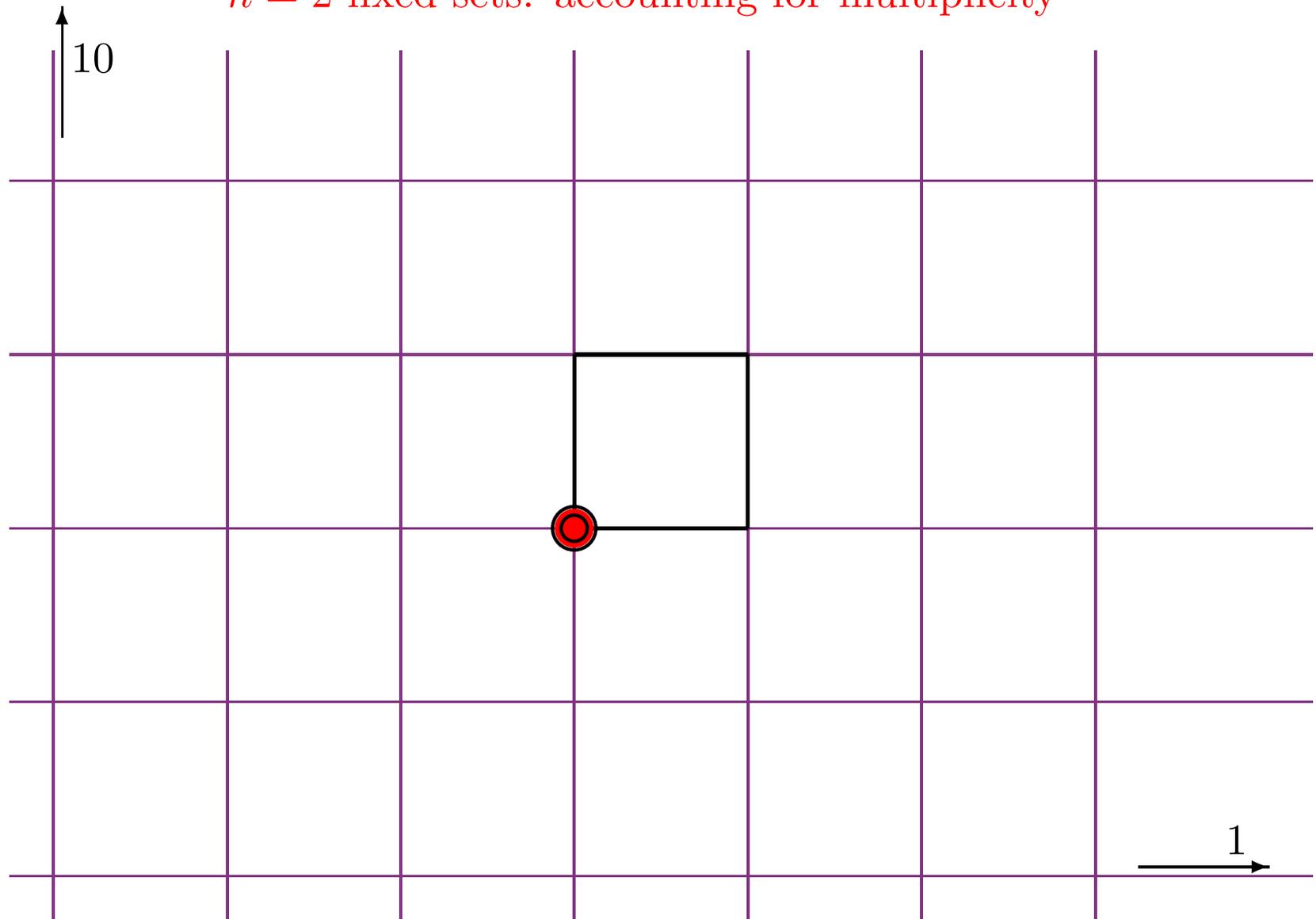
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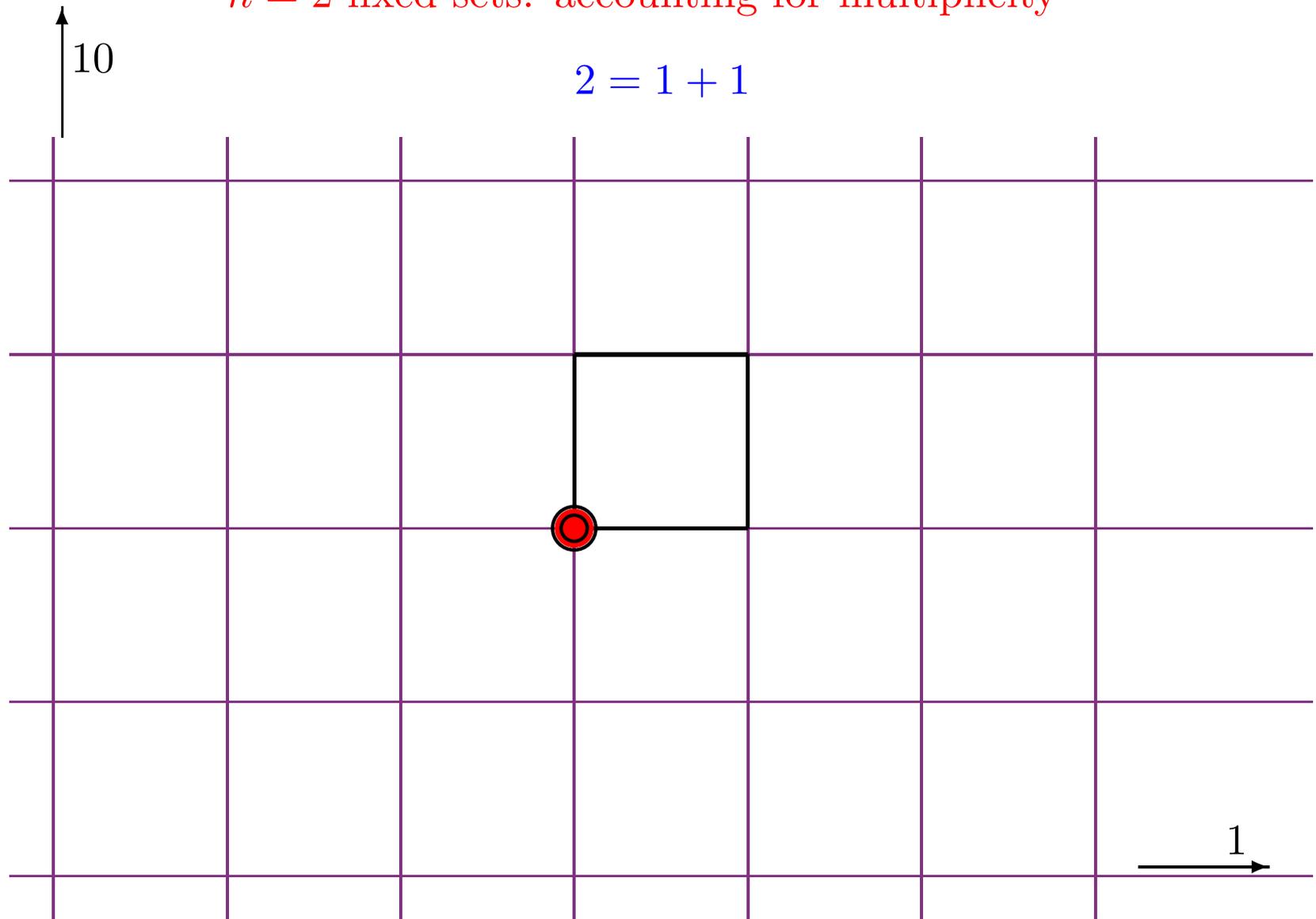


$n = 2$  fixed sets: accounting for multiplicity



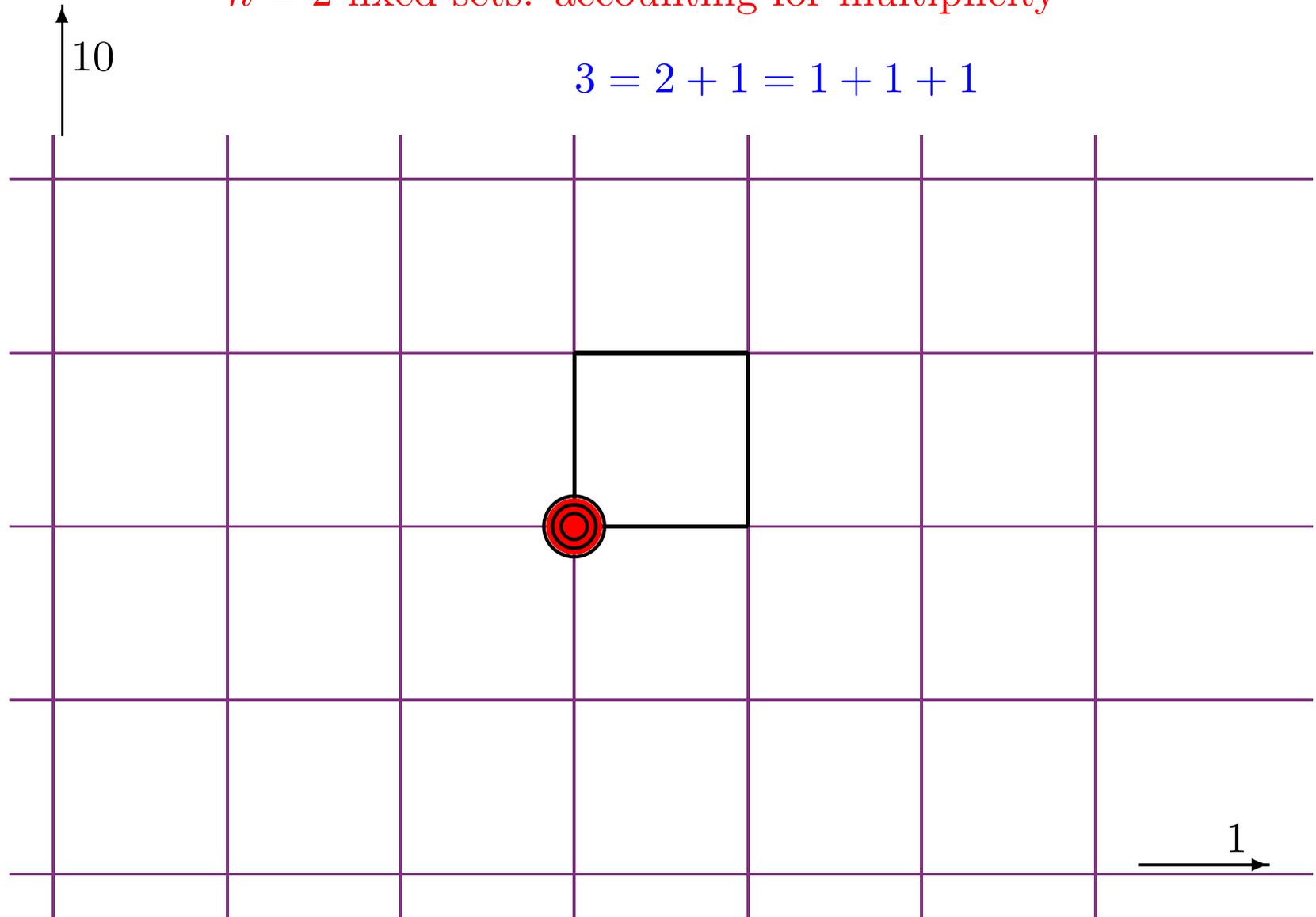
$n = 2$  fixed sets: accounting for multiplicity

$$2 = 1 + 1$$



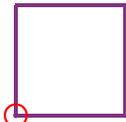
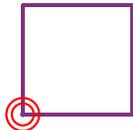
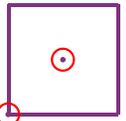
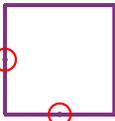
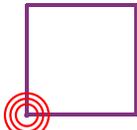
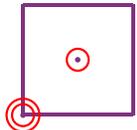
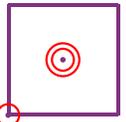
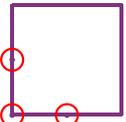
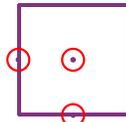
$n = 2$  fixed sets: accounting for multiplicity

$$3 = 2 + 1 = 1 + 1 + 1$$



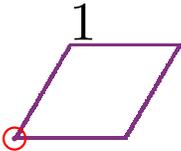
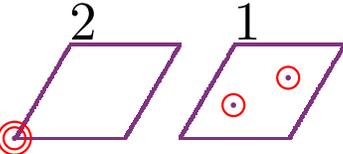
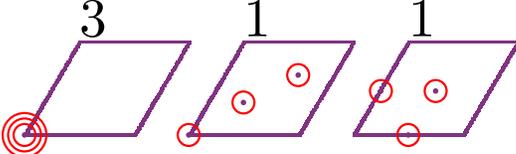
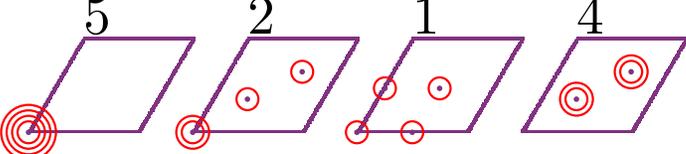
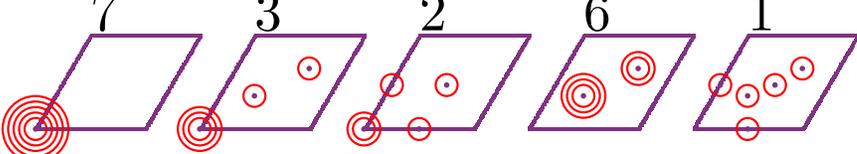
# Counting number of ground states

(Singlet RR ground state)

$v = \frac{\pi}{2}$	$n = 1$	<div style="display: flex; justify-content: space-around;"> <div style="text-align: center;">1 </div> <div style="text-align: center;">1 </div> </div>	$N_s = 2$
	$n = 2$	<div style="display: flex; justify-content: space-around;"> <div style="text-align: center;">2 </div> <div style="text-align: center;">2 </div> <div style="text-align: center;">1 </div> <div style="text-align: center;">1 </div> </div>	$N_s = 6$
	$n = 3$	<div style="display: flex; justify-content: space-around;"> <div style="text-align: center;">3 </div> <div style="text-align: center;">2 </div> <div style="text-align: center;">2 </div> <div style="text-align: center;">3 </div> <div style="text-align: center;">1 </div> <div style="text-align: center;">1 </div> </div>	$N_s = 12$

# Counting number of ground states

(Singlet RR ground state)

$\nu = \frac{\pi}{3}$	$n = 1$		$N_s = 1$
	$n = 2$		$N_s = 3$
	$n = 3$		$N_s = 5$
	$n = 4$		$N_s = 12$
	$n = 5$		$N_s = 19$

Number of states for  $U(n)$  on  $T^2$

$\tau$	$\nu$	$ k $	$n$				
			1	2	3	4	5
$e^{i\pi/3}$	$\frac{\pi}{3}$	1	<b>1</b>	<b>3</b>	<b>5</b>	<b>12</b>	<b>19</b>
$i$	$\frac{\pi}{2}$	2	<b>2</b>	<b>6</b>	<b>12</b>		
$e^{i\pi/3}$	$-\frac{2\pi}{3}$	3	<b>3</b>	<b>9</b>			

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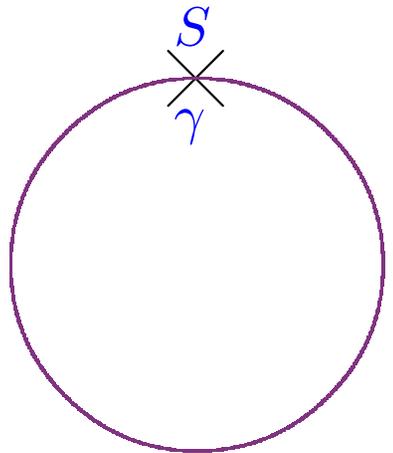
Chern-Simons:

$U(1)$  level  $k$ :  $N_s = k$ .

$SU(2)$  level  $k$ :  $N_s = k + 1$ .

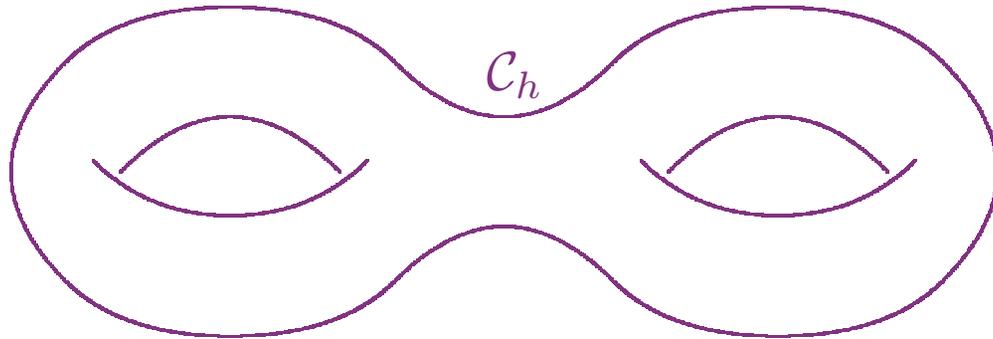
$SU(3)$  level  $k$ :  $N_s = (k + 1)(k + 2)/2$ .

S-twist  
& R-twist



$S^1$   
 $0 \leq x_3 < 2\pi R$

$\times$



Riemann surface  
area  $\mathcal{A}$

$\mathcal{A} \rightarrow 0 \implies \sigma\text{-model on } \mathcal{M}_H$

S-duality becomes T-duality [Bershadsky & Johansen & Sadv & Vafa; Harvey & Moore & Strominger]

# Witten Index

$$\#\{\text{vacua of 2+1D theory on } \mathcal{C}_h\} = I = \text{tr}_0\{(-1)^F \mathcal{T}(\mathbf{s})\gamma\}.$$

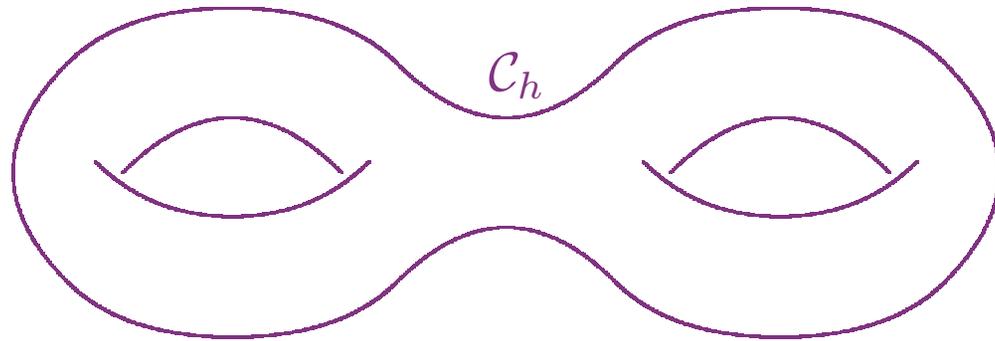
# Hitchin's equations

$$F_{z\bar{z}} = [\phi_z, \bar{\phi}_{\bar{z}}]$$

$$D_z \bar{\phi}_{\bar{z}} = D_{\bar{z}} \phi_z = 0$$

$A_z$  gauge field

$\phi_z$  adj.-valued 1-form



Riemann surface

# Hitchin's equations

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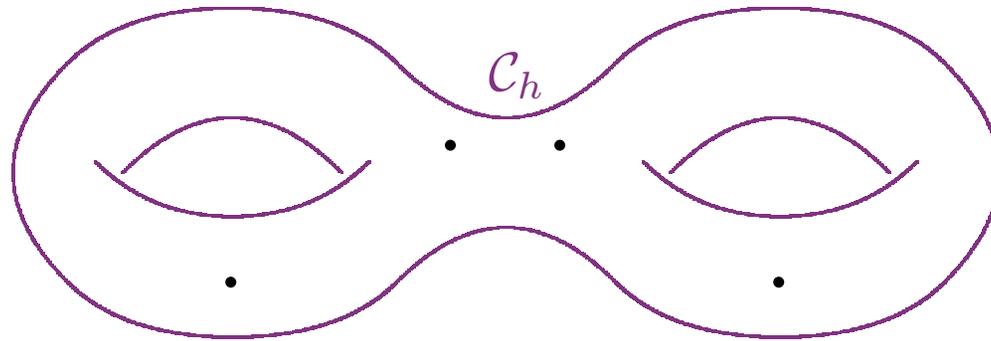
$$D_z \bar{\phi}_{\bar{z}} = D_{\bar{z}} \phi_z = 0$$

$b_{zz} = \text{tr}(\phi_z^2)$  holomorphic with  $4h - 4$  zeroes.

Space of quadratic differentials:  $\mathbb{C}^{3(h-1)}$

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Riemann surface

# Hitchin's equations

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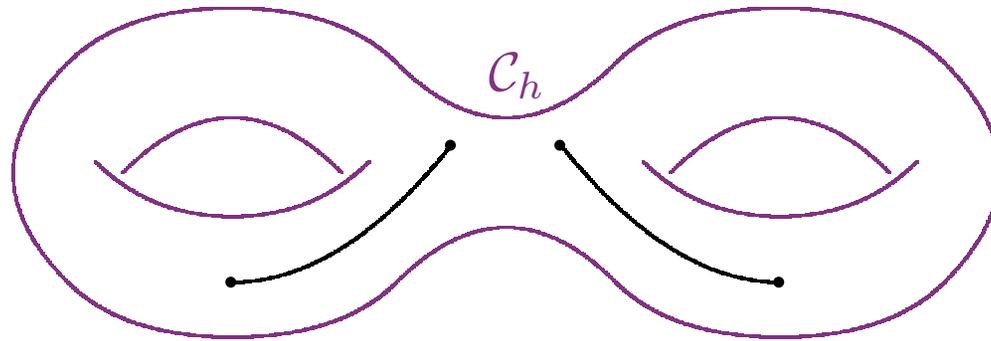
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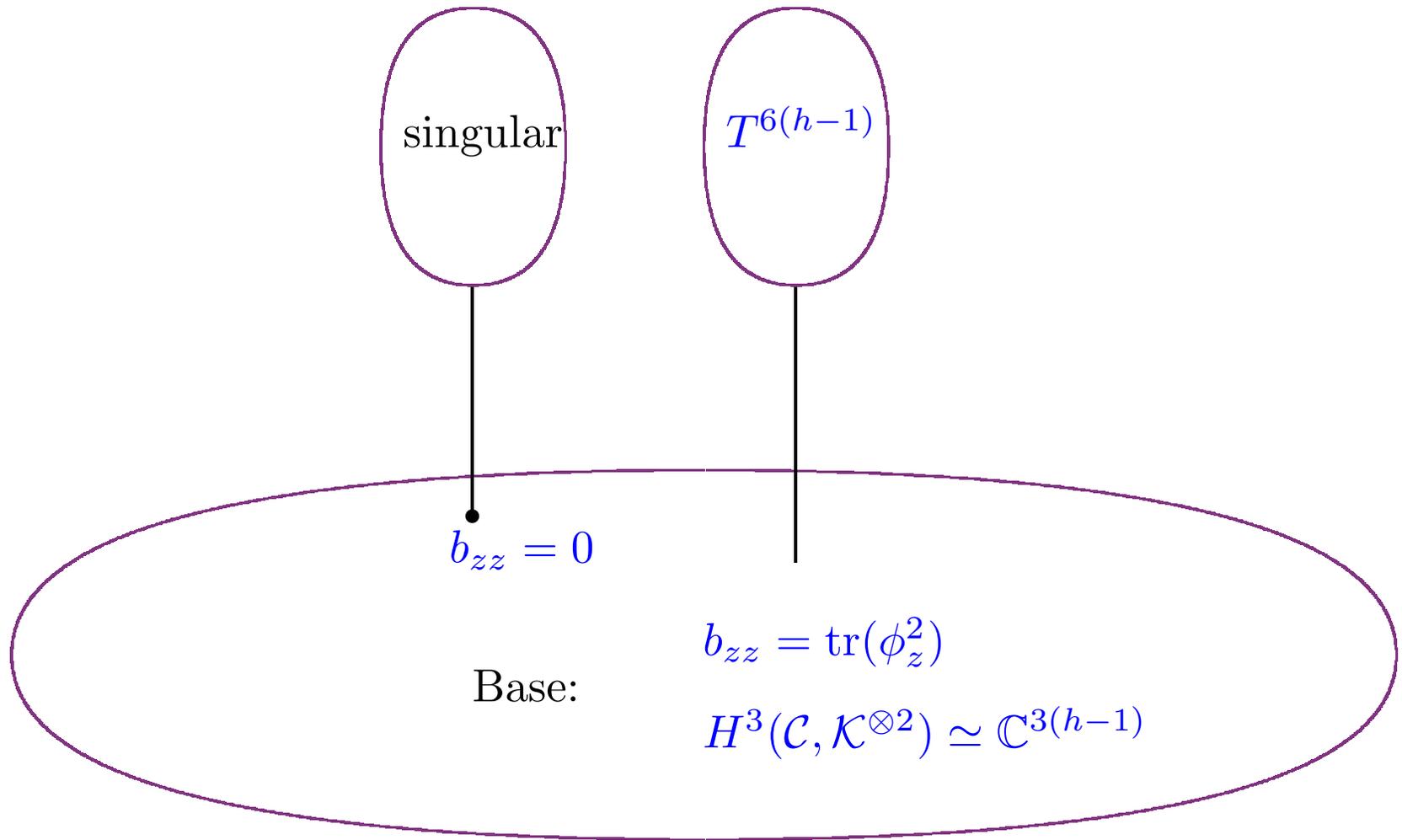


Riemann surface

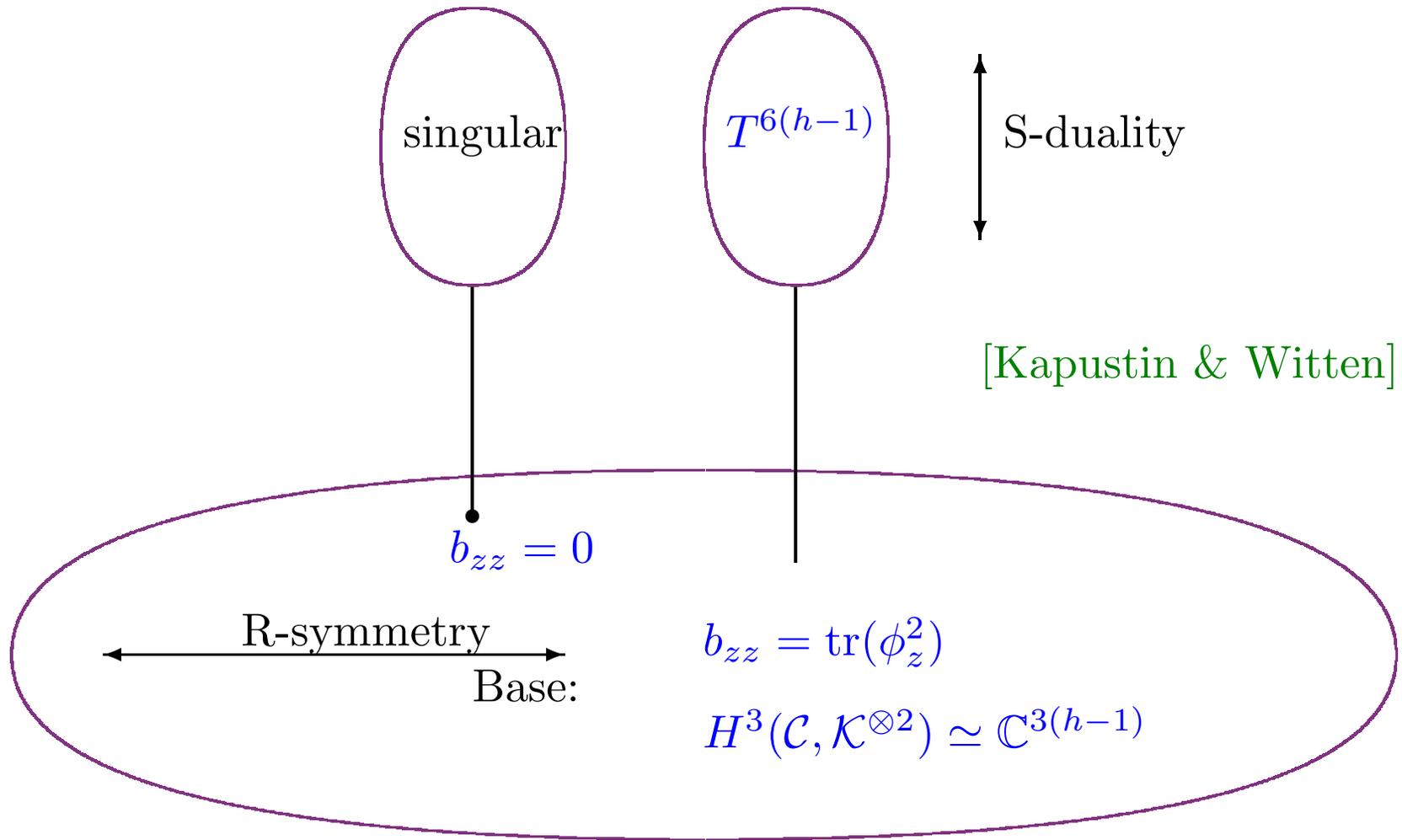
Double cover has genus  $4h - 3$

Prym subspace of its Jacobian:  $T^{6(h-1)}$

# Hitchin Fibration



# Hitchin Fibration



## The fiber over $b_{zz} = 0 \dots$

$$b_{zz} = \text{tr}(\phi_z^2) = 0$$

Case 1:  $\phi_z = 0 \implies \mathcal{M}_{\text{fc}} =$  moduli space of flat connections.

Case 2:  $\phi_z = \begin{pmatrix} 0 & \alpha_z \\ 0 & 0 \end{pmatrix}, \quad A_{\bar{z}} = \begin{pmatrix} a_{\bar{z}} & c_{\bar{z}} \\ 0 & -a_{\bar{z}} \end{pmatrix},$

$$a_{\bar{z}} = -\frac{1}{2} \partial_{\bar{z}} \log \alpha_z,$$

$$\partial_z a_{\bar{z}} - \partial_{\bar{z}} a_z = |\alpha_z|^2 + |c_{\bar{z}}|^2, \quad \text{and } \frac{c_{\bar{z}}^*}{\alpha_z} = \text{holomorphic.}$$

Special subcase of 2:  $c_{\bar{z}} = 0.$

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Case 1:  $\phi_z = 0 \implies \mathcal{M}_{\text{fc}}$  = moduli space of flat connections.

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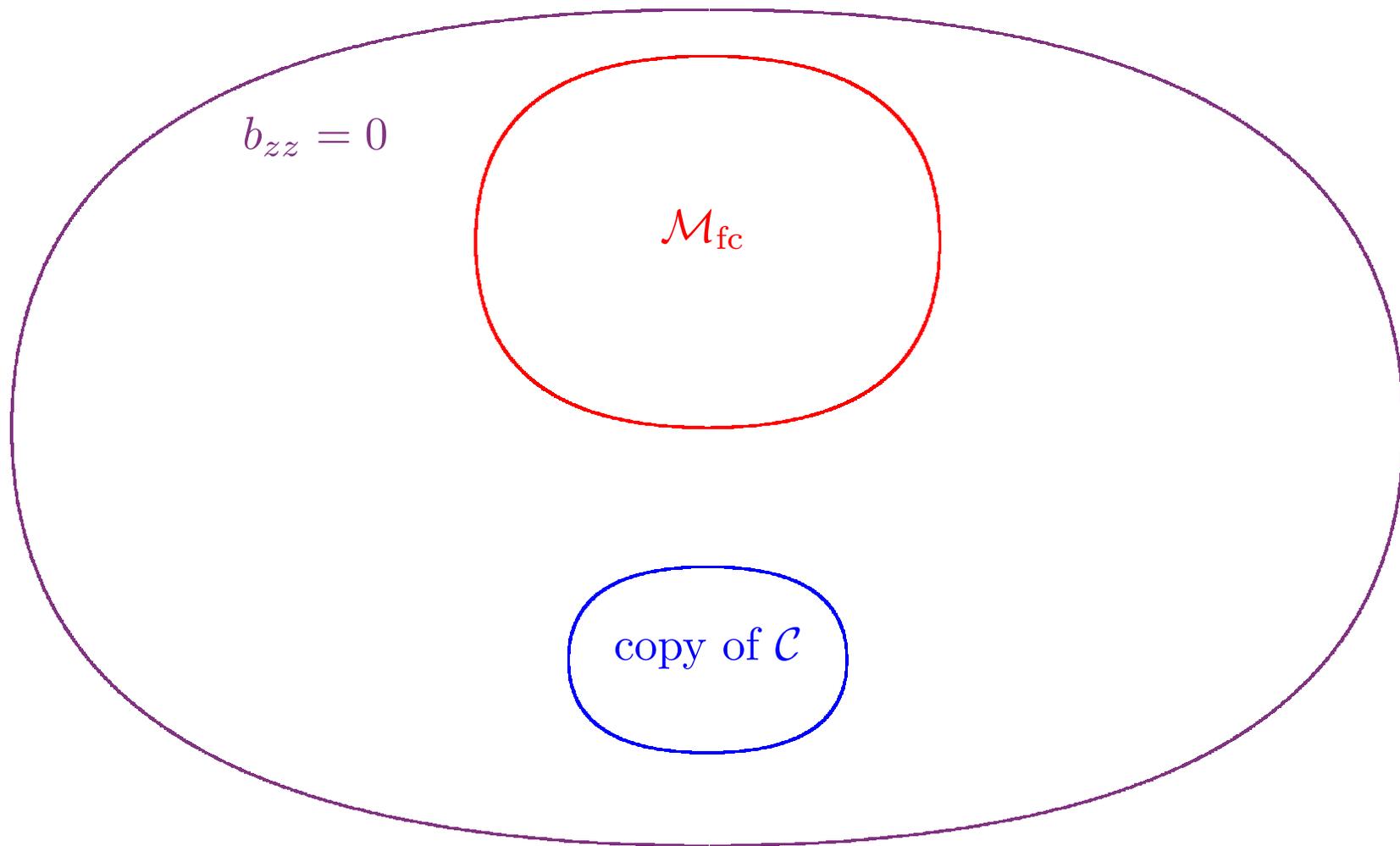
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Special subcase of 2:  $c_{\bar{z}} = 0$ .

if also genus  $h = 2$ :  $\alpha_z$  has a single simple zero on  $\mathcal{C}_2$  which determines the solution uniquely up to gauge.

The fiber over  $b_{zz} = 0 \dots$

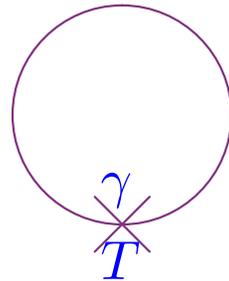


# T-duality and Geometric Quantization

1+1D  $\sigma$ -model with target space  $X$

$T =$  T-duality (mirror symmetry) twist

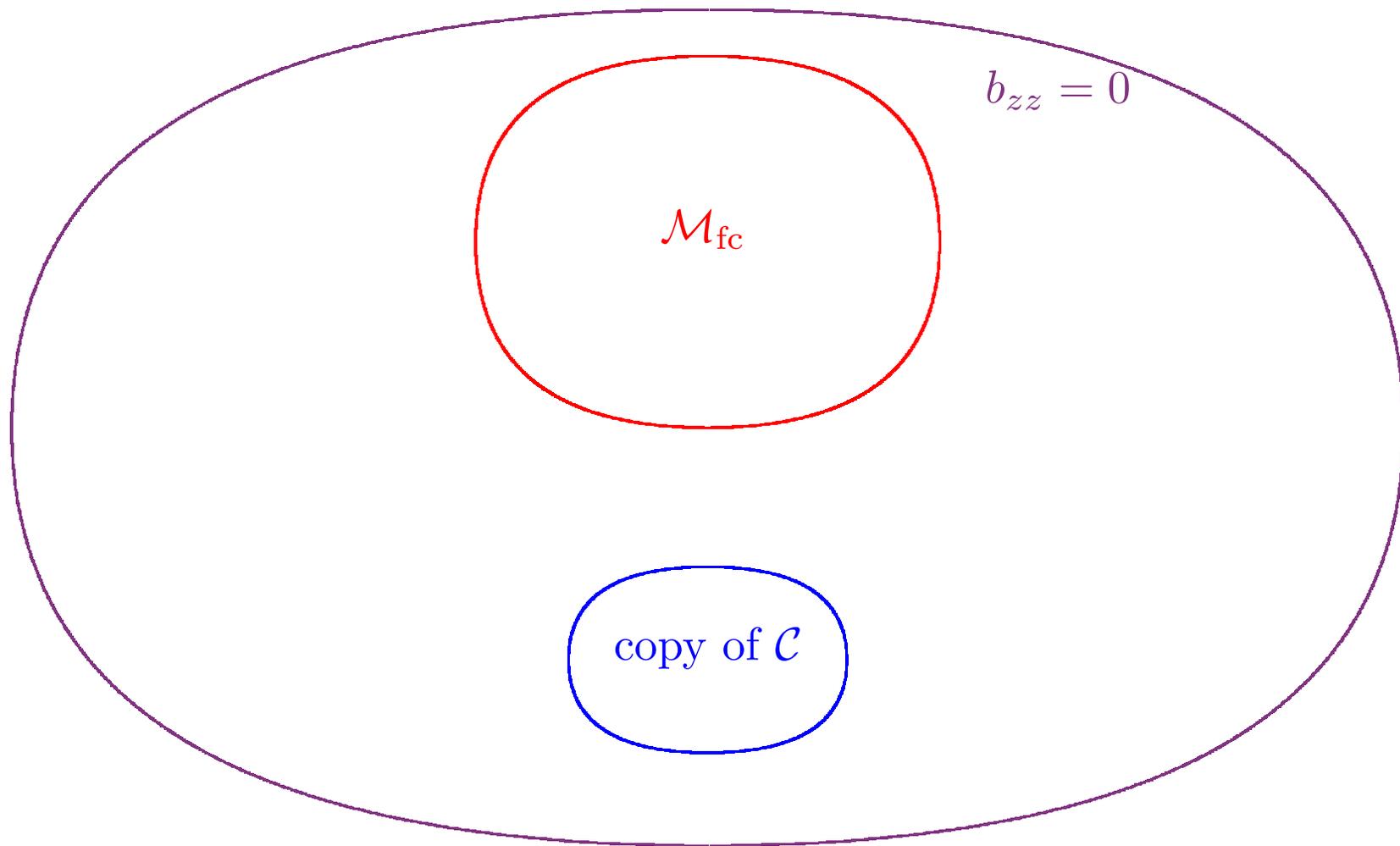
$\gamma =$  some isometry twist



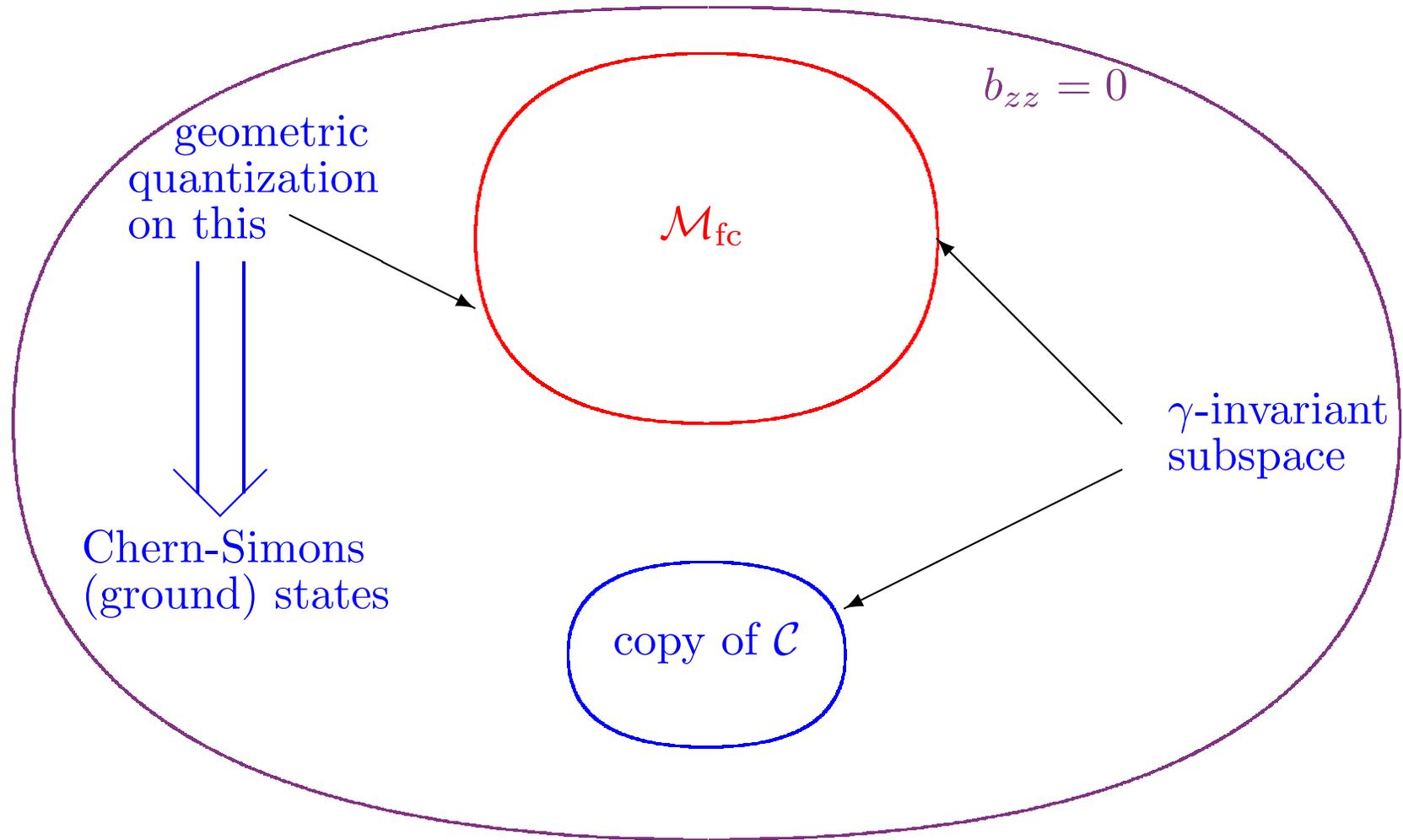
IR?

Geometric quantization on  $\gamma$ -invariant subspace???

The fiber over  $b_{zz} = 0 \dots\dots$



# The fiber over $b_{zz} = 0$ . . . . .



# Conclusions

- Compactification of  $N = 4 U(n)$  SYM on  $S^1$  with an S-duality twist, at a self-dual  $\tau$  seems to give a topological 2+1D QFT in IR for  $n$  sufficiently small;
- Number of (ground) states on  $T^2$  can be computed by string dualities;
- Number of (ground) states on  $\mathcal{C}_h$  ( $h > 1$ ) could be computed if we could determine the signs in the action of S-duality on  $H^*(\mathcal{M}_H)$ ;

## Open questions

- What is this topological 2+1D theory?
- Wilson lines?
- Mirror symmetry twist and geometric quantization?
- Nonlocal topological structure from the kernel  $\mathcal{S}(A, A_D)$ ?  
(Simple argument suggests that correlation functions of pairs of Wilson lines is proportional to the linking number.)
- Can we extract any new clues about S-duality from this?

**Thank you!**

**Title**